

A Linear Operator Inequalities/LMI Approach to Stability and Control of Distributed Parameter Systems

Emilia FRIDMAN*

Electrical Engineering, Tel Aviv University, Israel

April 3, 2011

Plan

- 1 Introduction
- 2 Exponential stability of linear time-delay systems in the Hilbert space
 - Linear Operator Inequalities for Exp. Stability in a Hilbert Space
 - LMIs for the Delay Heat Equation
 - LMIs for the Delay Wave Equation
 - Conclusions
- 3 Bounds on the response of a drilling pipe model
 - Motivation: *improved drilling towards Leviathan gas discovery*
 - Drilling system model: a 1-d wave equation
 - Ultimate boundedness

- Two main approaches are usually used for stability and control of infinite-dimensional systems:
 - ① the analysis and control of the abstract infinite-dimensional system (e.g. in the Hilbert space) with the corresponding conclusions for specific systems;
 - ② the direct approach to a specific system.
- In this talk both approaches to Lyapunov-based analysis will be presented:
 - ① the **Linear Operator Inequalities** (LOIs) for the stability of linear **time-delay systems in a Hilbert space**.
[Fridman & Orlov, Aut09a]
Thanks to *J.-P. RICHARD* for our visits to Ecole Centrale de Lille.
 - ② the **direct Lyapunov approach** to analysis of **1-d wave eq.**
[Fridman & Orlov, Aut09b], [Fridman, S. Mondie, B. Saldivar , IMA J.]

Plan

- 1 Introduction
- 2 Exponential stability of linear time-delay systems in the Hilbert space
 - Linear Operator Inequalities for Exp. Stability in a Hilbert Space
 - LMIs for the Delay Heat Equation
 - LMIs for the Delay Wave Equation
 - Conclusions
- 3 Bounds on the response of a drilling pipe model
 - Motivation: *improved drilling towards Leviathan gas discovery*
 - Drilling system model: a 1-d wave equation
 - Ultimate boundedness

- Delays may be a source of **instability**. However, they may have also a **stabilizing effect**.
- In the case of **distributed parameter systems**, arbitrarily **small delays** in the feedback may **destabilize** the system [**Datko**, SICON 88], [**Logemann et al.**, SICON 96], [**Nicaise & Pignotti**, SICON 06].
- Thus, the wave eq. non-robust w.r.t. delay [**Wang, Guo & Krstic**, SICON 11]:

$$\begin{aligned} z_{tt}(\xi, t) &= z_{\xi\xi}(\xi, t), & \xi \in (0, 1), \\ z(0, t) &= 0, & z_{\xi}(1, t) = kz_t(1, t - h) \end{aligned}$$

- is **stable** for $h = 0$ and $k = 1$ (all solutions are zero for $t \geq 2$), **unstable** for all **small enough** h and $k = 1$ [**Datko**, TAC 97]
- **stable** for $h = 2$ iff $k \in (0, 1)$, **unstable for arbitrary small perturbations of $h = 2$,**
- for $h = 2, 4, 6, 8$ (even multiples of the wave propagation) **stable** for some $k > 0$,
- for $h = 1, 3, 5, 0.5$ **unstable $\forall k$.**

- The stability analysis of PDEs with delay is essentially more complicated than of ODEs.
- There are only a few works on Lyapunov-based technique for PDEs with delay. The 2nd Lyapunov method was extended to abstract nonlinear time-delay systems in the Banach spaces in Wang (1994a, JMAA) and applied to [scalar heat and scalar wave equations](#) with constant delays and with the Dirichlet boundary conditions in Wang (1994b, JMAA), Wang(2006, JMAA).
- Stability and instability conditions for [wave delay](#) equations were found in (Nicaise & Pignotti, SIAM 2006).

- In (E. Fridman & Y. Orlov, Aut 09) *exp. stability of general* distributed parameter systems are derived for linear systems, where a *bounded operator acts on the delayed state*. The system delay is admitted to be *unknown and time-varying*.
- Sufficient *exp. stability* conditions are derived in the form of **Linear Operator Inequalities (LOIs)**, where the decision variables are operators in the Hilbert space.
General methods for solving LOI have not been developed yet. Some finite dimensional approximations were considered in Ikeda, Azuma & Uchida (2001).
- Being applied to a *heat/wave* equation these conditions are represented in terms of standard **finite-dimensional LMIs** that guarantee the stability of the 1-st/2-nd order delay-differential eqs. This *reduction of LOIs to finite-dimensional LMIs is tight*: the stability of the latter delay-differential eqs is **necessary** for the stability of the heat/wave eqs.

Problem Statement



$$\dot{x}(t) = Ax(t) + A_1x(t - \tau(t)), \quad t \geq t_0 \quad (1)$$

where $x(t) \in \mathcal{H}$, \mathcal{H} is a Hilbert space,
 delay $\tau(t)$ is a piecewise continuous function

$$\inf_t \tau(t) > 0, \quad \sup_t \tau(t) \leq h, \quad h > 0 \quad (2)$$

A_1 is a linear bounded operator,

A is an infinitesimal operator, generating a strongly continuous semigroup $T(t)$, the domain $\mathcal{D}(A)$ is dense in \mathcal{H} .

- Throughout, solutions of such a system are defined in the Caratheodory sense: (1) is required to hold almost everywhere.
- Let the initial conditions $x^{t_0} = \varphi(\theta)$, $\theta \in [-h, 0]$, $\varphi \in W$ be given in the space $W = C([-h, 0], \mathcal{D}(A)) \cap C^1([-h, 0], \mathcal{H})$.

Let the initial conditions

$$x^{t_0} \stackrel{\Delta}{=} x(t_0 + \theta) = \varphi(\theta), \theta \in [-h, 0], \varphi \in W$$

be given in the space $W = C([-h, 0], \mathcal{D}(A)) \cap C^1([-h, 0], \mathcal{H})$.

- Under the assumption

$$\inf_t \tau(t) = h_0 > 0, \sup_t \tau(t) \leq h, h > 0 \quad (3)$$

we have $\tau(t) \in [h_0, h]$.

The above initial-value problem is *well-posed* on $[t_0, \infty)$ and its solutions can be found as *mild solutions* of

$$\begin{aligned} x(t) &= T(t - t_0)x(t_0) \\ &+ \int_{t_0}^t T(t - s)A_1x(s - \tau(s))ds, \quad t \geq t_0. \end{aligned} \quad (4)$$

- In this talk we will consider 2 main examples:
 - 1) **heat** (parabolic eq); 2) **wave** (hyperbolic eq).
- **Example 1:** heat eq.

$$z_t(\xi, t) = az_{\xi\xi}(\xi, t) - a_1z(\xi, t - \tau(t)), \quad t \geq t_0, \quad 0 \leq \xi \leq \pi \quad (5)$$

with constants $a > 0$ and a_1 and with the Dirichlet b. c.

$$z(0, t) = z(\pi, t) = 0, \quad t \geq t_0. \quad (6)$$

a is the **heat conduction coefficient**,

a_1 is the **coefficient of the heat exchange** with the surroundings

$z(\xi, t)$ is the **temperature of the rod**

The above system describes the **propagation of heat in a homogeneous 1-d rod with a fixed temperature at the ends** in the case of the **delayed (possibly, due to actuation) heat exchange**.

- Heat eq. can be rewritten as

$$\dot{x}(t) = Ax(t) + A_1x(t - \tau(t)), \quad t \geq t_0$$

$\mathcal{H} = L_2(0, \pi)$, $A = a \frac{\partial^2}{\partial \xi^2}$ with the dense domain

$$\mathcal{D}\left(\frac{\partial^2}{\partial \xi^2}\right) = \{z \in W^{2,2}([0, \pi], \mathbf{R}) : z(0) = z(\pi) = 0\},$$

and with the **bounded operator** $A_1 = -a_1$.

- A **generates a strongly continuous semigroup** (see, e.g., Curtain & Zwart (1995) for details).

Example 2

:

- Wave equation

$$\begin{aligned} z_{tt}(\xi, t) &= az_{\xi\xi} - \mu_0 z_t(\xi, t) - a_0 z(\xi, t) \\ &\quad - a_1 z(\xi, t - \tau(t)), \quad t \geq 0, \quad 0 \leq \xi \leq \pi, \\ z(0, t) &= z(\pi, t) = 0, \quad t \geq t_0. \end{aligned} \quad (7)$$

Eqs (7) describe the **oscillations of a homogeneous string with fixed ends** in the case of the *delayed stiffness restoration*.

- Introduce the operators

$$A = \begin{bmatrix} 0 & 1 \\ a \frac{\partial^2}{\partial \xi^2} - a_0 & -\mu_0 \end{bmatrix}, \quad A_1 = \begin{bmatrix} 0 & 0 \\ -a_1 & 0 \end{bmatrix}$$

$$\mathcal{D}\left(\frac{\partial^2}{\partial \xi^2}\right) = \{z \in W^{2,2}([0, \pi], \mathbf{R}) : z(0) = z(\pi) = 0\},$$

- Then (7) can be represented as

$$\dot{x}(t) = Ax(t) + A_1x(t - \tau(t)), \quad t \geq t_0$$

in $\mathcal{H} = L_2(0, \pi) \times L_2(0, \pi)$ with the infinitesimal operator A ,

- possessing the domain $\mathcal{D}(A) = \mathcal{D}\left(\frac{\partial^2}{\partial \xi^2}\right) \times L_2(0, \pi)$ and generating a strongly continuous semigroup (see, e.g., Curtain & Zwart (1995)).

Our aim is to derive **exp. stability** criteria for (1), (2). The stability concept under study is based on the initial data norm in the space

$$W = C([-h, 0], \mathcal{D}(A)) \cap C^1([-h, 0], \mathcal{H})$$

defined as

$$\|\phi\|_W = |A\phi(0)| + \|\phi\|_{C^1([-h, 0], \mathcal{H})} \quad (8)$$

Suppose $x(t, t_0, \phi)$, $t \geq t_0$ denotes a solution of (1) with $x^{t_0} = \phi$. System (1) is said to be exponentially stable with a decay rate $\delta > 0$ if $\exists K \geq 1$:

$$|x(t, t_0, \phi)|^2 \leq Ke^{-2\delta(t-t_0)} \|\phi\|_W^2 \quad \forall t \geq t_0. \quad (9)$$

- Given a continuous functional $\mathbf{V} : \mathbf{R} \times \mathbf{W} \times \mathbf{C}([-h, 0], \mathcal{H}) \rightarrow \mathbf{R}$, \dot{V} along (1) is defined as follows:

$$\dot{V}(t, \phi, \dot{\phi}) = \limsup_{s \rightarrow 0^+} \frac{1}{s} [V(t+s, x^{t+s}(t, \phi), \dot{x}^{t+s}(t, \phi)) - V(t, \phi, \dot{\phi})].$$

- Lemma** Let $\exists \delta, \beta, \gamma$ and a continuous functional

$$V : \mathbf{R} \times W \times C([-h, 0], \mathcal{H}) \rightarrow \mathbf{R}$$

such that the function $\bar{V}(t) = V(t, x^t, \dot{x}^t)$ is absolutely continuous for x^t , satisfying (1), and

$$\begin{aligned} \beta |\phi(0)|^2 &\leq V(t, \phi, \dot{\phi}) \leq \gamma \|\phi\|_W^2, \\ \dot{V}(t, \phi, \dot{\phi}) + 2\delta V(t, \phi, \dot{\phi}) &\leq 0. \end{aligned} \tag{10}$$

Then (1) is exp. stable with the decay rate δ and with $K = \frac{\gamma}{\beta}$.

Notation:

Given a linear operator $\Phi : \mathcal{H} \rightarrow \mathcal{H}$ with a dense domain $\mathcal{D}(\Phi) \subset \mathcal{H}$, the notation Φ^* stands for the adjoint operator. Such an operator Φ is strictly positive definite, i.e., $\Phi > 0$, iff it is self-adjoint, i.e. $\Phi = \Phi^*$ and $\exists \beta > 0$ such that

$$\langle x, \Phi x \rangle \geq \beta \langle x, x \rangle, \forall x \in \mathcal{D}(\Phi)$$

$\Phi \geq 0$ means that $\langle x, \Phi x \rangle \geq 0$ for all $x \in \mathcal{D}(\Phi)$.

In a Hilbert space $\mathcal{D}(A)$, consider

$$V(t, x^t, \dot{x}^t) = \langle x(t), Px(t) \rangle + \int_{t-h}^t e^{2\delta(s-t)} \langle x(s), Sx(s) \rangle ds \\ + h \int_{-h}^0 \int_{t+\theta}^t e^{2\delta(s-t)} \langle \dot{x}(s), R\dot{x}(s) \rangle ds d\theta + \int_{t-\tau(t)}^t e^{2\delta(s-t)} \langle x(s), Qx(s) \rangle ds$$

$\mathbf{P} : \mathcal{D}(A) \rightarrow \mathcal{H}$ is a linear operator, $P > 0$,
 $R, Q, S \in \mathcal{L}(\mathcal{H})$, $R, Q, S \geq 0$
 $\forall x \in D(A)$ and some positive $\gamma_P, \gamma_Q, \gamma_S, \gamma_R$

$$\begin{aligned} \langle x, Px \rangle &\leq \gamma_P [\langle x, x \rangle + \langle Ax, Ax \rangle], & \langle x, Qx \rangle &\leq \gamma_Q \langle x, x \rangle, \\ \langle x, Rx \rangle &\leq \gamma_R \langle x, x \rangle, & \langle x, Sx \rangle &\leq \gamma_S \langle x, x \rangle \end{aligned} \quad (11)$$

By using **Cauchy-Schwartz (Jensen's) inequality**, we obtain conditions in 2 forms:

- 1) by substituting the right side of (1) for $\dot{x}(t)$;
- 2) by using **descriptor approach (Fridman SCL 2001)**:

Theorem 1 (1) is exp. stable with the decay rate δ if LOI is feasible

$$\Phi_h = \begin{bmatrix} \Phi_{11} & 0 & PA_1 \\ 0 & 0 & 0 \\ A_1^*P & 0 & 0 \end{bmatrix} + h^2 \begin{bmatrix} A^*RA & 0 & A^*RA_1 \\ 0 & 0 & 0 \\ A_1^*RA & 0 & A_1^*RA_1 \end{bmatrix} - e^{-2\delta h} \begin{bmatrix} R & 0 & -R \\ 0 & (S+R) & -R \\ -R & -R & 2R+(1-d)Q \end{bmatrix} \leq 0,$$

where $\Phi_h : \mathcal{D}(A) \times \mathcal{D}(A) \times \mathcal{D}(A) \rightarrow \mathcal{H} \times \mathcal{H} \times \mathcal{H}$ and where

$$\Phi_{11} = A^*P + PA + 2\delta P + Q + S. \quad (12)$$

Differently from the finite dimensional case, the feasibility of the strict LOIs for $h = 0$ (or $\delta = 0$) does not necessarily imply the feasibility of these LOIs for small enough h (δ) because h^2 (δ) is multiplied by the operator, which may be unbounded.

Theorem 1 gives **delay-dependent** conditions (h -dependent) even for $\delta \rightarrow 0$. For $S = R = 0$ we obtain the following "**quasi delay-independent**" conditions, which coincide for ODE with (Mondie & Kharitonov, TAC 2005):

Corollary Given $\delta > 0$, (1) is exp. stable with the decay rate δ for all delays with $\dot{\tau}(t) \leq d < 1$ if $\exists P > 0$ and $Q \geq 0$ subject to (11) such that the LOI

$$\begin{bmatrix} (A + \delta)^* P + P(A + \delta) + Q & PA_1 \\ A_1^* P & -(1 - d)Qe^{-2\delta h} \end{bmatrix} \leq 0 \quad (13)$$

holds in $\mathcal{D}(A) \times \mathcal{D}(A) \rightarrow \mathcal{H} \times \mathcal{H}$. The inequality (9) is satisfied with $K = \max\{\gamma_P, h\gamma_Q\} / \beta$.

It may be difficult to verify the feasibility of LOI 1, if **the operator that multiplies h^2 (and depends on A) in Φ_h is unbounded**. To avoid this, we will derive the *2-nd form* of LOI by the *descriptor method* (Fridman, SCL 2001), where the right-hand sides of the expressions

$$\begin{aligned} 0 &= 2\langle x(t), P_2^* [Ax(t) + A_1x(t - \tau(t)) - \dot{x}(t)] \rangle, \\ 0 &= 2\langle \dot{x}(t), P_3^* [Ax(t) + A_1x(t - \tau(t)) - \dot{x}(t)] \rangle \end{aligned} \quad (14)$$

with some $P_2, P_3 \in \mathcal{L}(\mathcal{H})$ are added into the right-hand side of \dot{V} .

LOI 2 via descriptor method

$$\begin{bmatrix} \Phi_{d11} & \Phi_{d12} & 0 & P_2^* A_1 + Re^{-2\delta h} \\ * & \Phi_{d22} & 0 & P_3^* A_1 \\ * & * & -(S+R)e^{-2\delta h} & Re^{-2\delta h} \\ * & * & * & -[2R + (1-d)Q]e^{-2\delta h} \end{bmatrix} \leq 0 \quad (15)$$

holds, where

$$\begin{aligned} \Phi_{d11} &= A^* P_2 + P_2^* A + 2\delta P + Q + S - Re^{-2\delta h}, \\ \Phi_{d12} &= P - P_2^* + A^* P_3, \\ \Phi_{d22} &= -P_3 - P_3^* + h^2 R. \end{aligned} \quad (16)$$

$$\begin{aligned} z_t(\xi, t) &= az_{\xi\xi}(\xi, t) - a_0z(\xi, t) - a_1z(\xi, t - \tau(t)), \\ t &\geq t_0, \quad 0 \leq \xi \leq \pi \end{aligned} \quad (17)$$

with constant $a > 0$, a_0, a_1 and with the Dirichlet boundary conditions

$$z(0, t) = z(\pi, t) = 0, \quad t \geq t_0. \quad (18)$$

Here we **apply the descriptor method LOIs**. The boundary-value problem (17), (18) can be rewritten as (1) in the Hilbert space

$\mathcal{H} = L_2(0, \pi)$ with $A = a \frac{\partial^2}{\partial \xi^2} - a_0$ with the dense domain

$$\mathcal{D}\left(\frac{\partial^2}{\partial \xi^2}\right) = \{z \in W^{2,2}([0, \pi], \mathbf{R}) : z(0) = z(\pi) = 0\}, \quad (19)$$

and with the bounded operator $A_1 = -a_1$.

A generates a strongly continuous semigroup

Delay-independent conditions are derived by using

$$V = p \int_0^\pi z^2(\xi, t) d\xi + q \int_{t-\tau(t)}^t \int_0^\pi e^{2\delta(s-t)} z^2(\xi, s) d\xi ds \quad (20)$$

with some constants $p > 0$ and $q > 0$. Then $P = p$, $Q = q$.

Integrating by parts and taking into account boundary conditions, we find

$$\langle x, (A^*P + PA)x \rangle = 2a \int_0^\pi p z z_{\xi\xi} d\xi = -2a \int_0^\pi p z_\xi^2 d\xi \leq -2a \int_0^\pi p z^2 d\xi$$

for $x \in \mathcal{D}(A)$, where the last inequality follows from the

Wirtinger's Inequality [Hardy et al., 88]: Let $z \in W^{1,2}([a, b], R)$ be a scalar function with $z(a) = z(b) = 0$. Then

$$\int_a^b z^2(\xi) d\xi \leq \frac{(b-a)^2}{\pi^2} \int_a^b (z'(\xi))^2 d\xi. \quad (21)$$

We thus obtain that the LOI is satisfied provided that the following LMI

$$\begin{bmatrix} q - 2(a + a_0)p & -a_1p \\ -a_1p & -(1 - d)qe^{-2\delta h} \end{bmatrix} < 0 \quad (22)$$

is feasible.

LMI (22) with $\delta = 0$ is feasible iff $\mathbf{a} + \mathbf{a}_0 > \mathbf{0}$, $\mathbf{a}_1^2 < (\mathbf{a} + \mathbf{a}_0)^2(1 - \mathbf{d})$.

Delay-Dependent Conditions

Choose the LKF of the form

$$\begin{aligned}
 V(t, z^t, z_s^t) &= (p_1 - p_3 a) \int_0^\pi z^2(\xi, t) d\xi + p_3 a \int_0^\pi z_\xi^2(\xi, t) d\xi \\
 &+ \int_0^\pi \left[r \int_{-h}^0 \int_{t+\theta}^t e^{2\delta(s-t)} z_s^2(\xi, s) ds d\theta \right. \\
 &\left. + s \int_{t-h}^t e^{2\delta(s-t)} z^2(\xi, s) ds + q \int_{t-\tau(t)}^t e^{2\delta(s-t)} z^2(\xi, s) ds \right] d\xi
 \end{aligned}$$

with $p_1 > 0, p_3 > 0, s > 0, r > 0$ and $q \geq 0$.

Then the operators in LOI take the form

$$P = -p_3 \left(a \frac{\partial^2}{\partial \xi^2} + a \right) + p_1, \quad R = r, \quad Q = q, \quad S = s,$$

$$P_3 = p_3, \quad P_2 = p_2 > 0, \quad p_2 - \delta p_3 \geq 0$$

Integrating by parts and applying Wirtinger's Inequality

$$\int_a^b z^2(\xi) d\xi \leq \frac{(b-a)^2}{\pi^2} \int_a^b (z'(\xi))^2 d\xi \quad (23)$$

we show that $P > 0$ and that the LOI of Th.2 is feasible if the following LMI is feasible

$$\begin{bmatrix} \phi_{11} & \phi_{12} & 0 & \phi_{14} \\ * & -2p_3 + h^2r & 0 & -p_3a_1 \\ * & * & -(s+r)e^{-2\delta h} & re^{-2\delta h} \\ * & * & * & \phi_{44} \end{bmatrix} < 0,$$

$$\phi_{11} = -2(a + a_0)p_2 + 2\delta p_1 + q + s - re^{-2\delta h},$$

$$\phi_{12} = p_1 - p_2 - (a + a_0)p_3, \quad \phi_{14} = -p_2a_1 + re^{-2\delta h},$$

$$\phi_{44} = -[2r + (1-d)q]e^{-2\delta h}$$

Remark

The same LMIs guarantee the exp. stability of the scalar equation

$$\dot{y}(t) + (a + a_0)y(t) + a_1y(t - \tau(t)) = 0 \quad (24)$$

Eq. (24) corresponds to **the first modal dynamics** (with $k = 1$) in the modal representation of the Dirichlet b. v. problem for the heat equation

$$y_k(t) + (a + a_0)k^2y_k(t) + a_1y_k(t - \tau(t)) = 0,$$

$k = 1, 2, \dots$ projected on the eigenfunctions of the operator $\frac{\partial^2}{\partial \xi^2}$ (this operator has eigenvalues $-k^2$).

The stability of the heat eq. implies the stability of ODE (24).

Thus the **reduction of infinite-dimensional LOI to finite-dimensional LMIs is tight: the stability of (24) is necessary for the stability of the heat eq.**

Remark

- The above LOIs (LMIs) are **affine in the system operators (coefficients)**. Consider now the systems under question with the uncertain operators (coefficients) from the **uncertain polytope, given by M vertices**. By the arguments of (Boyd *et al.* 1994), the uncertain systems are exp. stable if the corresponding **LOIs (LMIs) in the vertices are feasible**.

Example Consider the controlled heat eq.

$$\begin{aligned} z_t(\xi, t) &= z_{\xi\xi}(\xi, t) + rz(\xi, t) + u, \\ z(0, t) &= z(l, t) = 0, \end{aligned} \tag{25}$$

where $\xi \in (0, l)$, $t > 0$ and where r is uncertain parameter satisfying $|r| \leq \beta$.

- It was shown in (Rebiai & Zinober, IJC 1993) that for $l = 1$ $u = -\gamma z(\xi, t)$, $\gamma > (\frac{\beta}{2\pi})^2$ as. stabilizes (25).
By **our method** $u = -\gamma z(\xi, t)$, $\gamma > \beta - \pi^2$ Since $\beta - \pi^2 \leq (\frac{\beta}{2\pi})^2$, our method guarantees exp. stabilization via a **lower gain**.
- Consider next $l = \pi$, $\beta = 0.1$ and the feedback $u = -z(\xi, t - \tau(t))$ with the uncertain delay $\tau(t) \in [0, h]$, $\dot{\tau} \leq d < 1$
This is a **polytopic system reached by choosing** $r = \pm 0.1$. We verify the feasibility of delay-dependent LMIs in **2 vertices**: $r = \pm 0.1$. We use LMI toolbox of Matlab and find the maximum values of h for which the system remains **as. stable**:

$$d = 0.5, h = 2.04; \text{ unknown } d, h = 1.34.$$

The latter results correspond also to the stability of

$$\dot{y} = (-1 + r)y(t) - y(t - \tau(t)), \quad |r| \leq 0.1$$

Consider the wave equation

$$\begin{aligned} z_{tt}(\xi, t) = & az_{\xi\xi} - \mu_0 z_t(\xi, t) - \mu_1 z_t(\xi, t - \tau(t)) \\ & - a_0 z(\xi, t) - a_1 z(\xi, t - \tau(t)), \quad t \geq 0, 0 \leq \xi \leq \pi \end{aligned} \quad (26)$$

with the Dirichlet boundary condition (18).

Introduce the operators

$$A = \begin{bmatrix} 0 & 1 \\ a \frac{\partial^2}{\partial \xi^2} - a_0 & -\mu_0 \end{bmatrix}, \quad A_1 = \begin{bmatrix} 0 & 0 \\ -a_1 & -\mu_1 \end{bmatrix} \quad (27)$$

where the domain $\mathcal{D}(\frac{\partial^2}{\partial \xi^2})$ is determined by (19). Then (18), (26) can be represented as (1) in the Hilbert space $\mathcal{H} = L_2(0, \pi) \times L_2(0, \pi)$ with the infinitesimal operator A , possessing the domain

$\mathcal{D}(A) = \mathcal{D}(\frac{\partial^2}{\partial \xi^2}) \times L_2(0, \pi)$ and generating a strongly continuous semigroup (see, e.g., Curtain & Zwart (1995)).

Delay-Dependent Conds: $\mu_1 = 0$ We apply LOI of Theorem 1. Since the delay appears only in u , we choose V as follows:

$$\begin{aligned}
 V &= ap_3 \int_0^\pi z_\xi^2(\xi, t) d\xi + \int_0^\pi [z(\xi, t) \ z_t(\xi, t)] P_0 \begin{bmatrix} z(\xi, t) \\ z_t(\xi, t) \end{bmatrix} d\xi \\
 &+ \int_0^\pi \left[hr \int_{-h}^0 \int_{t+\theta}^t z_t^2(\xi, s) e^{2\delta(s-t)} ds d\theta + s \int_{t-h}^t z^2(\xi, s) e^{2\delta(s-t)} ds \right. \\
 &\left. + q \int_{t-\tau}^t z^2(\xi, s) e^{2\delta(s-t)} ds \right] d\xi, \\
 P_0 &= \begin{bmatrix} p_1 & p_2 \\ p_2 & p_3 \end{bmatrix}, \quad P_w = \begin{bmatrix} ap_3 + p_1 & p_2 \\ p_2 & p_3 \end{bmatrix} > 0.
 \end{aligned}$$

where $r > 0$, $s > 0$, $q \geq 0$. Then the operators P, Q, R in LOI are given by

$$\begin{aligned}
 P &= \begin{bmatrix} -ap_3 \frac{\partial^2}{\partial \xi^2} + p_1 & p_2 \\ p_2 & p_3 \end{bmatrix} > 0, \quad Q = \begin{bmatrix} q & 0 \\ 0 & 0 \end{bmatrix} \geq 0, \\
 R &= \text{diag}\{r, 0\} \geq 0, \quad S = \text{diag}\{s, 0\} \geq 0.
 \end{aligned}$$

Denote

$$C_\delta = \begin{bmatrix} \delta & 1 \\ -a - a_0 & -\mu_0 + \delta \end{bmatrix}. \quad (28)$$

LOI is feasible if the following LMIs are satisfied:

$$p_2 \geq p_3 \delta, \quad \begin{bmatrix} \phi_w & 0 & P_w \begin{bmatrix} 0 \\ -a_1 \end{bmatrix} + \begin{bmatrix} re^{-2\delta h} \\ 0 \end{bmatrix} \\ * & -(s+r)e^{-2\delta h} & re^{-2\delta h} \\ * & * & -(2r + (1-d)q)e^{-2\delta h} \end{bmatrix} < 0, \quad (29)$$

where

$$\phi_w = C_\delta^T P_w + P_w C_\delta + \text{diag}\{q + s - re^{-2\delta h}, h^2 r\}.$$

The same LMIs appear to guarantee the stability of ODE with delay

$$\dot{\bar{z}}(t) = C_0 \bar{z}(t) + A_1 \bar{z}(t - \tau(t)), \quad \bar{z}(t) \in \mathbf{R}^2.$$

This ODE governs the first modal dynamics of the modal representation of the Dirichlet boundary-value problem.

Example

- Consider the controlled wave equation

$$z_{tt}(\zeta, t) = 0.1z_{\zeta\zeta}(\zeta, t) - 2z_t(\zeta, t) + u, \quad (30)$$

- with boundary condition (18),
 $t \geq t_0, 0 \leq \zeta \leq \pi, 0 \leq \tau \leq h, \dot{\tau} \leq d < 1$.
- Applying LMI to the open-loop system we find that (30) with $u = 0$ is **exp. stable** with the decay rate $\delta = 0.05$.
- Considering next a delayed feedback
 $u = -z(\zeta, t - \tau(t))$
- and verifying the LMI, we find that the closed-loop system is exp. stable with a greater decay rate $\delta = 0.8$ for all $0 \leq \tau(t) \leq 0.31$.

Delay-Ind. Conditions are derived for $\mu_1 \neq 0$.

- **Corollary from delay-ind. conds** In a particular case where $a = 1$, $a_0 = a_1 = 0$, the Dirichlet boundary-value problem for the wave eq. is exp. stable for all $\tau \leq d < 1$ if $\mu_1^2 < (1 - d)\mu_0^2$
- **Remark.** The condition $0 \leq \mu_1 < \mu_0$ for the stability of the wave eq. with *constant delay* and $a = 1$, $a_0 = a_1 = 0$ was obtained by *Nicaise & Pignotti (SIAM 2006)*, where it was shown that if $\mu_1 \geq \mu_0$, there exists a sequence of arbitrary small delays that destabilize the system.

Thanks to **Jack Hale** for encouragement on Partial Functional Dif. Eqs and to **Chris Byrnes** for discussions on *Output Regulation for Distributed Parameter Systems with Delays*

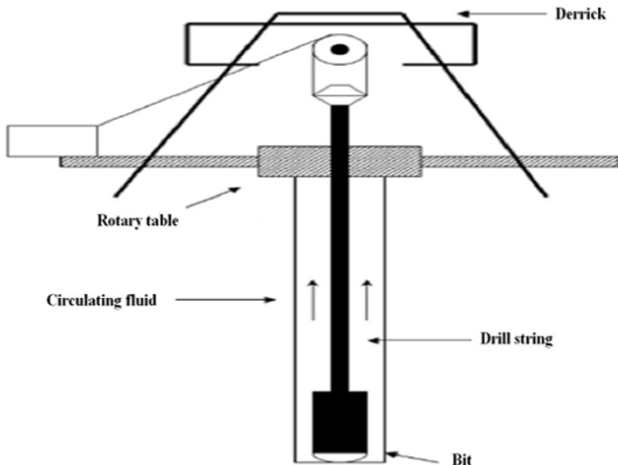
- A general framework is given for exp. stability of linear distributed parameter systems in a Hilbert space with a bounded operator acting on the delayed state.
- Stability conditions are derived in terms of LOIs in the Hilbert space.
- In the case of a heat/wave scalar equation with the Dirichlet boundary conditions, these LOIs are reduced to finite-dimensional LMIs by applying new Lyapunov functionals.
- The simplicity and the tightness of the results are the advantages of the new method.
- As it happened with LMIs in the finite dimensional case, LOIs are expected to provide effective tools for robust control of distributed parameter systems.

Plan

- 1 Introduction
- 2 Exponential stability of linear time-delay systems in the Hilbert space
 - Linear Operator Inequalities for Exp. Stability in a Hilbert Space
 - LMIs for the Delay Heat Equation
 - LMIs for the Delay Wave Equation
 - Conclusions
- 3 Bounds on the response of a drilling pipe model
 - Motivation: *improved drilling towards Leviathan gas discovery*
 - Drilling system model: a 1-d wave equation
 - Ultimate boundedness

Joint with Sabine Mondie and Belem Saldivar (Cinvestav, Mexico) (IMA J. 2010)

- ① Motivation: *improved drilling towards Leviathan gas discovery*
- ② Drilling system model
- ③ Ultimate boundedness:
 - difference equation approach
 - wave equation analysis
- ④ Numerical examples
- ⑤ Concluding remarks



A sketch of a simplified drillstring system is shown on Fig. 1.

Drilling system model

N. Challamel, *Rock destruction effect on the stability of a drilling structure*, Journal of sound and vibration, 233 (2), 235-254, 2000.

$$\frac{GJ}{L^2} \frac{\partial^2 z}{\partial \sigma^2}(\sigma, t) - I \frac{\partial^2 z}{\partial t^2}(\sigma, t) - \beta \frac{\partial z}{\partial t}(\sigma, t) = 0, \quad \sigma \in [0, 1],$$

$$z(0, t) = 0;$$

$$\frac{GJ}{L} \frac{\partial z}{\partial \sigma}(1, t) + I_B \frac{\partial^2 z}{\partial t^2}(1, t) = -T'(\Omega + \theta z_t) \frac{\partial z}{\partial t}(1, t) + w(t).$$

- $z(\sigma, t)$ is the angle of rotation,
- T is the torque on the bit,
- I_B is a lumped inertia (the assembly at the bottom hole),
- $\beta \geq 0$ damping (viscous and structural),
- I is the inertia, G is the shear modulus, J is the geometrical moment of inertia.

- The presence of a bounded additive noise signal $w(t)$ is considered at the bottom of the drillstring in order to account for external disturbances and modeling errors

$$|w(t)| \leq \bar{w}, \quad t \in (0, \infty).$$

- We will study **ultimate boundedness** of the solution.

The initial conditions are:

$$\begin{aligned} z(\sigma, 0) &= \zeta(\sigma), & z_\sigma(\sigma, 0) &= \dot{\zeta}(\sigma) \in L_2(0, 1), \\ z_t(\sigma, 0) &= \zeta_1(\sigma) \in L_2(0, 1). \end{aligned} \quad (31)$$

When the damping and the lumped inertia are negligible ($\beta = I_B = 0$)
the model reduces to:

$$\frac{\partial^2 z}{\partial t^2}(\sigma, t) = a \frac{\partial^2 z}{\partial \sigma^2}(\sigma, t), \quad \sigma \in [0, 1], t \geq 0 \quad (32)$$

$$z(0, t) = 0; \quad \frac{\partial z}{\partial \sigma}(1, t) = -k \frac{\partial z}{\partial t}(1, t) + r w(t) \quad (33)$$

where $a = \frac{GI}{IL^2}$, $k = \frac{LI'}{GJ}$, $r = \frac{L}{GJ} \in \mathbb{R}$.

A differential-difference equation approach

The general solution of the unidimensional wave equation is:

$$z(\sigma, t) = \phi(t + s\sigma) + \psi(t - s\sigma), \quad t \geq s,$$

where $\phi, \psi \in C^1$ and $s = \sqrt{\frac{1}{a}}$.

The boundary conditions can be rewritten as:

$$z(0, t) = \phi(t) + \psi(t) = 0,$$

$$\frac{\partial z(1, t)}{\partial \sigma} = s\dot{\phi}(t + s) - s\dot{\psi}(t - s) = -k[\dot{\phi}(t + s) + \dot{\psi}(t - s)] + rw(t).$$

A differential-difference equation approach

It follows from the above expressions that

$$\phi(t) = -\psi(t), \quad t \geq 0.$$

We obtain the **differential-difference equation** for $t \geq s$:

$$\dot{\psi}(t+s) = -\frac{(s-k)}{(s+k)}\dot{\psi}(t-s) - \frac{r}{(s+k)}w(t), \quad s = \sqrt{\frac{1}{a}}$$

with the initial condition

$$\begin{aligned} \dot{\psi}(t) &= -0.5[\zeta_1(t/s) + \check{\zeta}(t/s)/s], \quad t \in [0, s], \\ \dot{\psi}(-t) &= 0.5[\zeta_1(t/s) - \check{\zeta}(t/s)/s], \quad t \in [0, s]. \end{aligned}$$

A differential-difference equation approach

Notice that

$$\frac{\partial z}{\partial \sigma}(\sigma, t) = s\dot{\phi}(t + s\sigma) - s\dot{\psi}(t - s\sigma), \quad \frac{\partial z}{\partial t}(\sigma, t) = \dot{\phi}(t + s\sigma) + \dot{\psi}(t - s\sigma).$$

We now summarize the above results:

Lemma

The solution of the boundary-value problem (32), (33) is ultimately bounded and for all initial conditions (31) it satisfies the inequalities

$$|z_\sigma(1, t)| \leq se^{-\frac{\lambda}{s}t} [|\zeta_1(\xi/s)| + |\dot{\zeta}(\xi/s)| / s] + \frac{2s|c_1|}{1 - e^{-\lambda}} \bar{w}, \quad t \geq 0,$$

$$|z_t(1, t)| \leq e^{-\frac{\lambda}{s}t} [|\zeta_1(\xi/s)| + |\dot{\zeta}(\xi/s)| / s] + \frac{2|c_1|}{1 - e^{-\lambda}} \bar{w} \quad t \geq 0.$$

A wave equation analysis

Under the assumption that the lumped inertia is negligible (i.e. $I_B = 0$) the model reduces to:

$$z_{tt}(\sigma, t) = az_{\sigma\sigma}(\sigma, t) + dz_t(\sigma, t) \quad t \geq t_0, \quad 0 \leq \sigma \leq 1$$

with the boundary conditions:

$$\begin{aligned} z(0, t) &= 0, \\ z_\sigma(1, t) &= -kz_t(1, t) + rw(t), \quad t \geq 0, \end{aligned}$$

where $a = \frac{GJ}{IL^2}$, $d = \frac{-\beta}{I}$, $r = \frac{L}{GJ}$ and $k = \frac{LT'}{GJ}$ with $0 < k_0 \leq k \leq k_1$.

$$\begin{aligned} z(\sigma, 0) &= \zeta(\sigma), \quad z_\sigma(\sigma, 0) = \dot{\zeta}(\sigma) \in L_2(0, 1), \\ z_t(\sigma, 0) &= \zeta_1(\sigma) \in L_2(0, 1). \end{aligned}$$

A wave equation analysis

Outline of the proof.

We look for an energy function V such that:

Lemma (Fridman & Dambrine, 2008)

Assume that $|w| \leq \bar{w}$. Let $V : [0, \infty) \rightarrow \mathbb{R}^+$ be an absolutely continuous function. If there exists $\delta > 0$, $b > 0$ such that the derivative of V satisfies almost everywhere the inequality

$$\frac{d}{dt}V + 2\delta V - bw^2 \leq 0,$$

then it follows that

$$V(t) \leq e^{-2\delta(t-t_0)}V(t_0) + (1 - e^{-2\delta(t-t_0)})\frac{b}{2\delta}\bar{w}^2.$$

A wave equation analysis

Consider the **energy** function

$$V(z_\sigma(\cdot, t), z_t(\cdot, t)) = p \left[\int_0^1 a z_\sigma^2(\sigma, t) d\sigma + \int_0^1 z_t^2(\sigma, t) d\sigma \right] + 2\chi \int_0^1 \sigma z_\sigma(\sigma, t) z_t(\sigma, t) d\sigma$$

(Nicaise & Pignotti, 2006) with constants $p > 0$ and small enough χ .
Following **LMI approach of (Fridman & Y. Orlov, Aut 09)**

$$\begin{aligned} V(z_\xi(\cdot, t), z_t(\cdot, t)) &\geq \int_0^1 [z_\xi \ z_t] \begin{bmatrix} a_1 p & \chi \xi \\ \chi \xi & p \end{bmatrix} [z_\xi \ z_t]^T d\xi \\ &> \varepsilon \int_0^1 [z_\xi^2(\xi, t) + z_t^2(\xi, t)] d\xi \end{aligned} \quad (34)$$

for some $\varepsilon > 0$. The latter inequality holds if

$$\begin{bmatrix} a_1 p & \chi \xi \\ \chi \xi & p \end{bmatrix} > 0, \quad \forall \xi \in [0, 1] \Leftrightarrow \begin{bmatrix} a_1 p & \chi \\ \chi & p \end{bmatrix} > 0$$

A wave equation analysis

$$\begin{aligned}
 \frac{d}{dt}V &= 2p \int_0^1 az_\sigma(\sigma, t)z_{t\sigma}(\sigma, t)d\sigma + 2p \int_0^1 z_t(\sigma, t)z_{tt}(\sigma, t)d\xi \\
 &\quad + 2\chi \frac{d}{dt} \left(\int_0^1 \sigma z_t z_\sigma d\sigma \right) \\
 &= 2p \int_0^1 [az_\sigma(\sigma, t)z_{t\sigma}(\sigma, t) + az_t(\sigma, t)z_{\sigma\sigma}(\sigma, t)]d\sigma \\
 &\quad + 2pd \int_0^1 z_t^2(\sigma, t)d\sigma + 2\chi \frac{d}{dt} \left(\int_0^1 \sigma z_t z_\sigma d\sigma \right).
 \end{aligned}$$

Integration by parts + boundary conditions \Rightarrow

$$\begin{aligned}
 \int_0^1 z_t(\sigma, t)z_{\sigma\sigma}(\sigma, t)d\sigma &= z_t(\sigma, t)z_\sigma(\sigma, t)|_0^1 - \int_0^1 z_{t\sigma}(\sigma, t)z_\sigma(\sigma, t)d\sigma \\
 &= z_t(1, t)(-kz_t(1, t) + rw(t)) - \int_0^1 z_{t\sigma}(\sigma, t)z_\sigma(\sigma, t)d\sigma.
 \end{aligned}$$

$$\begin{aligned}
 \frac{d}{dt} \left(2 \int_0^1 \xi z_t z_\xi d\xi \right) &= - \int_0^1 (z_t^2 + az_\sigma^2)d\sigma + z_t^2(1, t) \\
 &\quad + a[-kz_t(1, t) + w(t)]^2 + 2d \int_0^1 \sigma z_t(\sigma, t)z_\sigma d\sigma.
 \end{aligned}$$

We find

$$\frac{d}{dt}V + 2\delta V - bw^2 = \int_0^1 \vartheta^T(\sigma, t) \Psi \vartheta(\sigma, t) d\sigma \leq 0$$

with

$$\vartheta^T(\sigma, t) = [z_t(1, t) \quad z_\sigma(\sigma, t) \quad z_t(\sigma, t), w(t)]$$

if

$$\Psi = \begin{bmatrix} \psi_1 & 0 & 0 & -a\chi kr + apr \\ * & \psi_2 & (2\delta + d)\chi\sigma & 0 \\ * & * & \psi_3 & 0 \\ * & * & * & -b + \chi ar^2 \end{bmatrix} < 0 \Leftrightarrow \Psi|_{\sigma=1} < 0$$

where

$$\psi_1 = -2akp + (1 + ak^2)\chi,$$

$$\psi_2 = -a\chi + 2\delta ap,$$

$$\psi_3 = -\chi + 2pd + 2\delta p.$$

A wave equation analysis

Theorem

Given $\delta > 0$, if $\exists p > 0, \chi > 0$ such that

$$\Psi|_{\sigma=1, k=k_i} < 0, \quad i = 0, 1, \quad \begin{bmatrix} a_1 p & \chi \\ \chi & p \end{bmatrix} > 0$$

then

$$\int_0^1 [z_\sigma^2(\sigma, t) + z_t^2(\sigma, t)] d\sigma \leq \frac{\alpha_2}{\alpha_1} e^{-2\delta(t-t_0)} \int_0^1 [\dot{\zeta}(\sigma)^2 + \zeta_1(\sigma)^2] d\sigma + \frac{b}{\alpha_1 2\delta} \bar{w}^2$$

with

$$\alpha_1 = \lambda_{\min} \begin{bmatrix} ap & 0 \\ 0 & p \end{bmatrix}, \quad \alpha_2 = \lambda_{\max} \begin{bmatrix} ap & \chi \\ \chi & p \end{bmatrix}.$$

is satisfied.

A numerical example

For the parameter values given in (Challamel, 1999) and $\beta = 0$:

$$\begin{aligned} G &= 79.3 \times 10^9 \text{ N/m}^2, & I &= 0.095 \text{ Kg} \cdot \text{m}, \\ T &= 3000 \text{ N} \cdot \text{m}, & J &= 1.19 \times 10^{-5} \text{ m}^4, \\ L &= 3145 \text{ m}, \end{aligned}$$

Difference equation approach \Rightarrow

$$\begin{aligned} |z_\sigma(\mathbf{1}, t)| &\leq 0.9979 e^{-0.2006t} [|\zeta_1(\xi/0.9979)| \\ &\quad + 1.0021 |\dot{\zeta}(\xi/0.9979)|] + 0.0033 \bar{w}. \end{aligned}$$

$$\begin{aligned} |z_t(\mathbf{1}, t)| &\leq e^{-0.2006t} [|\zeta_1(\xi/0.9979)| \\ &\quad + 1.0021 |\dot{\zeta}(\xi/0.9979)|] + 0.0033 \bar{w}. \end{aligned}$$

A numerical example: Wave equation approach

The wave equation approach leads to

Case	1	2	3	4	5
δ	0.08	0.06	0.04	0.01	0.0001
b	3.2521	1.0707	1.2145	1.5221	1.7951

For $\delta = 0.04$

$$\int_0^1 [z_\sigma^2(\sigma, t) + z_t^2(\sigma, t)] d\sigma \leq 1.1909e^{-0.08t} \int_0^1 [\zeta_1^2(\sigma) + \dot{\zeta}^2(\sigma)] d\sigma + 11.9944\bar{w}^2.$$

The *difference equation* approach leads to an ultimate bound for the main variable of interest, the *angular velocity at the drill bottom* $z_t(1, t)$, while the *wave equation* model provides the bound on the *energy*

$$\int_0^1 [z_\sigma^2(\sigma, t) + z_t^2(\sigma, t)] d\sigma.$$

Recent results

on *Robust Stability and Control of Distributed Parameter Systems*

- Boundary value H_∞ control of the heat/wave eqs
(with [Y. Orlov](#), Aut09).
Thanks to *J.-P. RICHARD* for our visits to Ecole Centrale de Lille.
- 2-nd order evolution eqs with *boundary tvr delays*
(with [S. Nicaise & J. Valein](#), SICON 10)
Thanks to *M. DAMBRINE* for my visits to Valenciennes Uni.
- *Sampled-data control* of 1-d semilinear heat equation
under the discrete in space and in time measurements
(with [A. Blighovsky](#), IFAC Congress 2011)

CONCLUSIONS

Many important plants (flexible manipulators and heat transfer processes) are governed by PDEs and described by uncertain models.

- The existing results [Bensoussan et al 1993], [Curtain & Zwart, 1995], [Foias et al. 1996] , [van Keulen 1993] on robust control of Distributed Parameter Systems (DPS) are devoted to the linear case.
- The LMI approach (Fridman & Orlov, Aut09a,b) is appropriate for nonlinear distributed parameter models and provides the desired system performance in spite of significant uncertainties.

As it happened with Time-Delay systems, LMIs are expected to provide effective tools for robust control of Distributed Parameter Systems.